Solid-State Electronics 67 (2012) 105-108



Contents lists available at ScienceDirect

Solid-State Electronics

journal homepage: www.elsevier.com/locate/sse



High-frequency, 6.2 Å pN heterojunction diodes

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ARTICLE INFO

Article history: Received 19 May 2011 Accepted 19 July 2011 Available online 8 September 2011

The review of this paper was arranged by Prof. A. Zaslavsky

Keywords: Semiconductor devices Diodes Mixers Frequency conversion

ABSTRACT

Sb-based pN heterojunction diodes at 6.2 Å, consisting of narrow bandgap p-type $In_{0.27}Ga_{0.73}Sb$ and wide bandgap n-type $In_{0.69}Al_{0.41}As_{0.41}Sb_{0.59}$, have been fabricated and measured. These diodes show excellent electrical characteristics with an ideality factor of 1.2 and high current density. S-parameter measurements and subsequent analysis show that these diodes have RC-cutoff frequencies over 1 THz, making these diodes excellent choices for high-frequency applications, such as sub-harmonic mixers for frequency conversion.

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1. Introduction

The 6.1 Å materials, as they are commonly referred to, InAs, AlSb, GaSb, and their alloys (e.g., In_{0.2}Al_{0.8}Sb, InAs_{0.9}Sb_{0.1}) have become highly desirable for use in low-power, high-speed electronic applications due to a large range of available bandgaps and band offsets and high electron and hole mobilities. High electron mobility transistors (HEMTs) fabricated from these materials have shown good operating characteristics [1,2]. Furthermore, the first monolithic microwave integrated circuits (MMICs) fabricated using 6.1 Å based HEMTs have been demonstrated [3]. New materials such as $In_xGa_{1-x}Sb$, $InAs_ySb_{1-y}$, and $In_xAl_{1-x}As_ySb_{1-y}$, with lattice constants ranging from 6.1 Å to 6.48 Å, show promise of further power reduction, due to narrower bandgaps, while maintaining or possibly improving high-speed operation [4]. These new materials' potential has been already demonstrated in HEMTs [5,6], diodes [7], and heterojunction bipolar transistors (HBTs) [8] at lattice constants around 6.2 Å. Additionally, increased interest in imaging technology in the sub-terahertz to terahertz (THz) frequency range [9-11], especially in limited power applications, has resulted in further examination of these materials for devices needed to produce high-frequency mixers and multipliers.

For frequency conversion, traditionally anti-parallel Schottky diode pairs have been used as sub-harmonic mixers. Schottky diodes have been preferable over pn diodes due to the latter's relatively higher overall capacitance and subsequently reduced frequency performance. The higher capacitance seen in pn diodes is explained

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as due to the addition of a diffusion capacitance associated with the diffusion of excess minority carriers within the neutral bulk of the device. Restriction to the use of Schottky diodes limits the fabrication of these devices to relatively larger bandgap materials (e.g., GaAs), as simple Schottky contacts to narrower bandgap materials (e.g., InAs) are practically impossible.

However, the work by Laux and Hess suggests that diffusion capacitance in pn diodes is not nearly as high as previously believed [12]. Additionally, pn diodes can effectively function as unipolar devices by "removing" one carrier type by means of a heterostructure barrier (e.g., $\Delta \mathcal{E}_v$ greater than a few kT blocks hole injection). The end result is a greatly reduced diffusion capacitance for the heterojunction pn diode and subsequently improved frequency performance. Therefore, heterojunction pn diodes fabricated from 6.1 to 6.3 Å materials are expected to exceed the performance of current state-of-the-art GaAs Schottky diodes.

In this letter, MMIC-compatible Sb-based pN heterojunction diodes at a lattice constant of a = 6.2 Å, consisting of narrow bandgap p-type In_{0.27}Ga_{0.73}Sb and wide bandgap n-type In_{0.69}Al_{0.41}As_{0.41}Sb_{0.59}, are presented that show excellent electrical behavior with RC-cutoff frequencies (f_{RC}) exceeding 1 THz.

2. Materials and methods

The pN diodes presented in this paper were grown by solid-source molecular beam epitaxy (MBE). Here, the use of a lower-case letter (p) for the narrow bandgap layer and upper-case letter (N) for the wide bandgap layer follows that established by Kroemer [13,14]. Beginning with a semi-insulating (SI) GaAs substrate, the growth buffer consisted of 1000 Å unintentionally doped (UID)

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1. REPORT DATE JUL 2011		2. REPORT TYPE		3. DATES COVE 00-00-201 1	red I to 00-00-2011	
4. TITLE AND SUBTITLE			5a. CONTRACT NUMBER			
High-frequency, 6.2	on diodes	5b. GRANT NUMBER				
				5c. PROGRAM E	LEMENT NUMBER	
6. AUTHOR(S)				5d. PROJECT NUMBER		
				5e. TASK NUMBER		
				5f. WORK UNIT NUMBER		
7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES) Naval Research Laboratory, Electronics Science and Technology Division, Washington, DC, 20375				8. PERFORMING ORGANIZATION REPORT NUMBER		
9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES)				10. SPONSOR/MONITOR'S ACRONYM(S)		
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12. DISTRIBUTION/AVAIL Approved for public	ABILITY STATEMENT	on unlimited				
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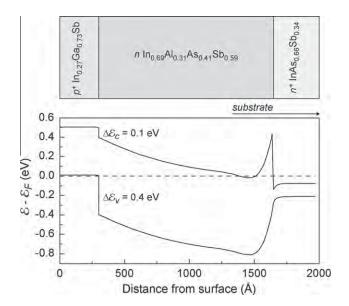


Fig. 1. Layer structure and band diagram of the pN diode examined in this paper. Bandgaps and band offsets are calculated following the work of Vurgaftman et al. [15]. The band diagram was simulated using ATLAS by SILVACO International.

GaAs (a=5.65 Å), 5000 Å UID Al_{0.65}Ga_{0.35}Sb (a=6.12 Å), and 1 µm UID In_{0.21}Ga_{0.19}Al_{0.60}Sb (a=6.2Å). After the buffer, the growth was continued with a 5000 Å n^+ (Te: 3×10^{18} cm⁻³) InAs_{0.66}Sb_{0.34} contact layer; a In_{0.69}Al_{0.31}As_{0.41}Sb_{0.59} wide bandgap n-layer consisting of an initial 50 Å n^+ (Te: 4×10^{18} cm⁻³) layer, a 250 Å n-type doping grade (Te: 4×10^{18} to 5×10^{16} cm⁻³), a 1000 Å n (Te: 5×10^{16} cm⁻³) "bulk" layer, and a 50 Å UID spacer; and a 300 Å p^+ (Be: 3×10^{19} cm⁻³) In_{0.27}Ga_{0.73}Sb layer (Fig. 1).

InAs_{0.66}Sb_{0.34} and In_{0.27}Ga_{0.73}Sb were selected for the *n*-type contact and *p*-type contact layers, respectively, due to the excellent transport ($\mu_{n,InAsSb}$ = 5600 cm²/Vs, $\mu_{p,InGaSb}$ = 160 cm²/Vs) and contact ($r_{c,n-InAsSb}$ = 2.4 × 10⁻⁸ Ω cm², $r_{c,p-InGaSb}$ = 8.9 × 10⁻⁸ Ω cm²) properties both have shown [16,17]. The relatively large valence band offset ($\Delta \mathcal{E}_{v} \approx 0.3$ eV) between In_{0.27}Ga_{0.73}Sb (\mathcal{E}_{g} = 0.49 eV) and In_{0.69}Al_{0.31}As_{0.41}Sb_{0.59} (\mathcal{E}_{g} = 0.83 eV) effectively makes the diode a unipolar electron device by blocking hole injection from the *p*-type In_{0.27}Ga_{0.73}Sb into the *n*-type In_{0.69}Al_{0.31}As_{0.41}Sb_{0.59}. Additionally, the use of a thin or "short" *p*-type layer reduces any capacitance associated with the diffusion of minority electrons across the layer.

Diode fabrication began with the definition and deposition of Pd:Pt:Au (100:50:2500 Å) unannealed, *p*-type ohmic contacts onto

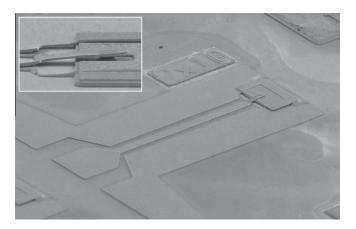


Fig. 2. Scanning electron microscope (SEM) image of a *pN* diode, including the coplanar wavequide (CPW). The inset shows a magnified view of the diode and bridge.

the p^+ In_{0.27}Ga_{0.73}Sb layer using standard e-beam lithography and e-beam evaporation. After which diode mesas were etched down to the InAs_{0.66}Sb_{0.34} contact layer by a SiCl₄-based reactive-ion etch (RIE), using lithographically-defined photoresist as the etch mask. Subsequently, Ti:Pt:Au (100:50:2500 Å) unannealed, n-type ohmic contacts were defined and deposited onto the n^+ InAs_{0.66}Sb_{0.34} contact layer. Diodes were isolated down to the SI GaAs substrate with a H₃PO₄-based wet etch. Diode fabrication was completed with a RF-compatible co-planar waveguide (CPW) and airbridge, between the diode and CPW (Fig. 2). Both circular and rectangular diodes were fabricated with sizes down to $A_{circ} = \pi (5/2)^2 \ \mu \text{m}^2 \approx 20 \ \mu \text{m}^2$ and $A_{rect} = 2 \times 5 \ \mu \text{m}^2 = 10 \ \mu \text{m}^2$, respectively.

3. Measurements and discussion

Two-point current–voltage (JV) measurements on the diodes show excellent behavior with an ideality factor of $\eta \approx 1.2$, strong rectifying behavior, and area-dependent scaling (Fig. 3). As can be seen in the figure, the diodes do not become resistively limited until $V \approx V_{bi} = 0.39$ V, suggesting that the series resistance of these diodes is very small.

One-port scattering parameter (S-parameter) measurements were taken of the diodes. The RF probes were calibrated on a separate calibration substrate, with open and short waveguide standards measured on-wafer. The waveguide was designed to be nominally 50 Ω for the range of frequencies measured (10 MHz–40 GHz). S_{11} of the as-measured device (waveguide and diode) and de-embedded diode for a 1 \times 10 μ m² area diode, at an applied voltage of 0 V, are shown in Fig. 4.

A small-signal model consisting of a parallel resistance (R_j) and capacitance (C_j) , representative of the diode's junction, in series with a second resistance (R_S) , representing any series resistance due to the neutral, undepleted bulk and contacts, was used in the extraction of the various device components (Fig. 5), with the total extracted junction capacitance being the combination of the depletion and diffusion capacitance $(C_j = C_{dep} + C_{diff})$.

Diodes in general are useful up to frequencies given by f_{RC} = 1/ $2\pi R_S C_j$. At frequencies above f_{RC} , the junction capacitance effectively appears as a short, resulting in the diode appearing as a simple resistance with a value of R_S and the characteristic current–voltage relationship of the diode is lost. The small-signal model was applied to both the as-measured device and de-embedded diode S-parameters of various area diodes. The resultant RC-cutoff frequencies, f_{RC} = 1/ $2\pi R_S C_j$, from the extracted component values are plotted in Fig. 6.

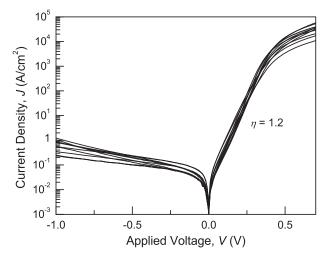


Fig. 3. Plot of the current–voltage (JV) measurements for circular and rectangular pN diodes of various area. The ideality factor for all measurements was $\eta \approx 1.2$.

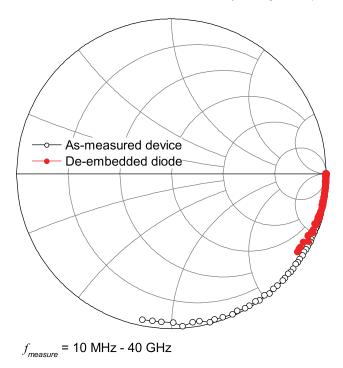


Fig. 4. S_{11} of the as-measured device (waveguide and diode, hollow symbols) and de-embedded diode (solid symbols) for a $1\times10~\mu\text{m}^2$ area diode, at an applied voltage of 0 V. The measurement frequency range was 10 MHz–40 GHz.

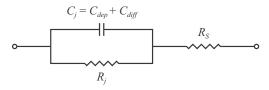


Fig. 5. Small-signal model of the pN diode. R_j and C_j represent the diode's junction resistance and capacitance, respectively, and R_S represents the total resistance outside of the junction due to the neutral, undepleted bulk and contacts.

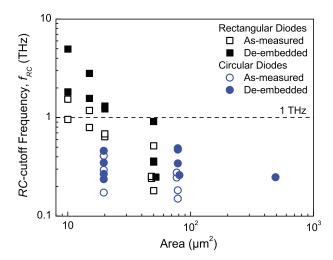


Fig. 6. Plot of the as-measured device and de-embedded diode *RC*-cutoff frequencies $(f_{RC}=1/2\pi R_S C_j)$ for circular and rectangular pN diodes of various area, calculated from the S-parameter extracted series resistance (R_S) and total junction capacitance (C_j) . De-embedded diodes with an area $\leq 20 \ \mu m^2$ exhibit $f_{RC} > 1$ THz. Note, overlapping points have been shifted for clarity.

As seen in the figure, even with the response of the waveguide included (hollow symbols), pN diodes with areas under $15~\mu m^2$ have RC-cutoff frequencies over 1 THz. For de-embedded pN diodes with areas under $20~\mu m^2$, RC-cutoff frequencies over 1 THz are observed (solid symbols, Fig. 6). Due to the limited frequency range used for the S-parameter measurements (10~MHz-40~GHz), accurate extraction of the the diode components, in particular R_S , for the smaller area diodes was difficult and requires higher frequency measurements to verify, but it is clear that f_{RC} exceeds 1 THz for these devices.

In general, rectangular diodes resulted in higher *RC*-cutoff frequencies, especially those with high L/W aspect ratios, as compared to circular diodes of similar area (Fig. 6). This is believed to be primarily due to the lower spreading resistance under the mesa for high L/W aspect ratio rectangular mesas ($R_{rect} = \rho_{sheet}/12 \times (L/W)^{-1}$, $R_{circ} = \rho_{sheet}/8\pi$)[18].

4. Conclusions

Sb-based pN heterojunction diodes at a = 6.2 Å, consisting of narrow bandgap, p-type $In_{0.27}Ga_{0.73}Sb$ and wide bandgap, n-type $In_{0.69}Al_{0.41}As_{0.41}Sb_{0.59}$, have been presented that show excellent electrical behavior ($\eta \approx 1.2$) with RC-cutoff frequencies (f_{RC}) exceeding 1 THz. This performance is a direct result of the excellent electronic properties (e.g., high mobility and low contact resistance) of $In_{0.27}Ga_{0.73}Sb$, $In_{0.69}Al_{0.31}As_{0.41}Sb_{0.59}$, and $InAs_{0.66}Sb_{0.34}$ and reduced capacitance associated with the pN heterojunction diode design.

Acknowledgement

This work was supported by the Office of Naval Research.

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